

A MODIFIED CONVOLUTION AND PRODUCT THEOREM FOR THE LINEAR CANONICAL TRANSFORM DERIVED BY REPRESENTATION TRANSFORMATION IN QUANTUM MECHANICS

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The Linear Canonical Transform (LCT) is a four parameter class of integral transform which plays an important role in many fields of signal processing. Well-known transforms such as the Fourier Transform (FT), the FRactional Fourier Transform (FRFT), and the FreSnel Transform (FST) can be seen as special cases of the linear canonical transform. Many properties of the LCT are currently known but the extension of FRFTs and FTs still needs more attention. This paper presents a modified convolution and product theorem in the LCT domain derived by a representation transformation in quantum mechanics, which seems a convenient and concise method. It is compared with the existing convolution theorem for the LCT and is found to be a better and befitting proposition. Further, an application of filtering is presented by using the derived results.

Keywords: linear canonical transform, convolution and product theorem, quantum mechanical representation.

1. Introduction

The Linear Canonical Transform (LCT) (Alieva and Bastiaans, 1999; Moshinsky and Quesne, 1971; Nazarathy and Shamir, 1982; Pei and Ding, 2001; 2002; Hennelly and Sheridan, 2005) is a four parameter (a, b, c, d) class of linear integral transform. It was introduced in the 1970s (Moshinsky and Quesne, 1971) and has found many applications in signal processing, radar system analysis, filter design, phase retrieval, pattern recognition (Shin and Park, 2011; Krzysztof and Domańska, 2011), signal synthesis, and many other areas (Stern, 2006; Tao *et al.*, 2004; Ozaktas *et al.*, 2000; Almeida, 1994; Barshan *et al.*, 1997; Sharma and Joshi, 2006; Goel and Singh, 2011). Well-known transforms such as the Fourier Transform (FT), the FRactional Fourier Transform (FRFT), and the FreSnel Transform (FST) can be seen as special cases of the linear canonical transform (Tao *et al.*, 2004; Ozaktas *et al.*, 2000; Almeida, 1994). Origins of LCTs are in quantum mechanics; a brief overview may be found in the work of Wolf (1979). As a generalization of the FT and FRFT, the

basic theories of the LCT have been developed including sampling the signals (Deng *et al.*, 2006; Li *et al.*, 2007; Koc *et al.*, 2008; Healy and Sheridan, 2008; Tao *et al.*, 2008), discrete approximations to the transforms (Pei and Ding, 2002; Hennelly and Sheridan, 2005; Healy and Sheridan, 2009; Oktem and Ozaktas, 2010), and so on, which can enrich the theoretical system of the LCT.

Many properties of the LCT and FRFT are already known (Alieva and Bastiaans, 1999; Pei and Ding, 2002; Ozaktas *et al.*, 2000), including the product and convolution theorem (Deng *et al.*, 2006; Ozaktas *et al.*, 1994; Almeida, 1997; Zayed, 1998; Sharma and Joshi, 2007; Wei *et al.*, 2012; 2009; Singh and Saxena, 2011), but none have received acclamation because either their definitions do not generalize very nicely to the classical definition for the FT or result in a bigger hardware complexity due to a large number of chirp signals.

In mathematics and, in particular, functional analysis, convolution is a mathematical operation on two functions f and g , producing a third function that is typically viewed as a modified version of one of the original functions, giving the area overlap between the

two functions as a function of the amount by which one of the original functions is translated. In other words, the Fourier transform of the convolution of two signals is the point-wise product of the Fourier transform of their respective signals. In the literature, many definitions of the convolution theorem for the LCT (Deng *et al.*, 2006; Wei *et al.*, 2012; 2009) are proposed. The classical definition of the convolution and product theorem of the FT for the signals $f(x)$ and $g(x)$ is given by

Convolution:

$$f(x) \otimes g(x) = \int_{-\infty}^{\infty} f(y)g(x-y) dy \xrightarrow{FT} \sqrt{2\pi} F(p) G(p), \quad (1)$$

Product:

$$f(x) \cdot g(x) \xrightarrow{FT} F(p) \otimes G(p), \quad (2)$$

where $F(p)$ and $G(p)$ are the FTs of $f(x)$ and $g(x)$, respectively, and ‘ \otimes ’ denotes the convolution operation.

In this paper, a modified convolution and product theorem is proposed with a new proof from the quantum mechanical representations transform, which seems simple and elegant. With the help of simulations, an attempt is made to compare the simulation results of the proposed theorem with the theorems proposed in the literature and it is found that the proposed theorem is a better proposition and can be considered a generalization of the classical convolution theorem of the FT. In the application section, a filter is designed by using the derived results.

2. Brief review of the correspondence between the LCT and its quantum mechanical counterpart

The LCT is also known as the generalized Fresnel transform (James and Agarwal, 1996; Palma and Bagini, 1997), the ABCD transform (Bernardo, 1996), the Collins formula (Collins, 1970), generalized Huygens integrals (Nazarathy and Shamir, 1982), quadratic-phase integrals or quadratic-phase systems (Bastiaans, 1979), extended fractional Fourier transforms (Hua *et al.*, 1997) as well as Moshinsky and Quesne integrals (Moshinsky and Quesne, 1971), and is a special case of the special affine Fourier transform (Abe and Sheridan, 1994a; 1994b). The one-dimensional LCT with parameters of a signal $f(x)$ is defined (Alieva and Bastiaans, 1999) as

$$L_{(a,b,c,d)}[f(x)](p) = \begin{cases} \int_{-\infty}^{+\infty} f(x)K_{(a,b,c,d)}(p,x)dx, & b \neq 0, \\ \sqrt{d}e^{j(cd/2)p^2} f(d \cdot p), & b = 0, \end{cases} \quad (3)$$

where parameters (a, b, c, d) are real numbers and the matrix $\begin{pmatrix} a & b \\ c & d \end{pmatrix}$ belongs to $SL(2, \mathbb{R})$.

In this paper, we only consider the case of $b \neq 0$ since the LCT is just a chirp multiplication operation if $b = 0$. The term $K_{(a,b,c,d)}(p, x)$ represents the integral kernel and is given by

$$K_{(a,b,c,d)}(p, x) = \sqrt{\frac{1}{j2\pi b}} \exp \left[\frac{j(ax^2 + dp^2)}{2b} - \frac{jxp}{b} \right]. \quad (4)$$

As special cases of the LCT, when $(a, b, c, d) = (\cos \alpha, \sin \alpha, -\sin \alpha, \cos \alpha)$, the LCT reduces to the FRFT, and when $(a, b, c, d) = (0, 1, -1, 0)$ the LCT reduces to the FT.

To relate the LCT kernel in a quantum mechanical representation, we use the completeness relation of the coherent state with the arguments of z_1 and z_2 . The coherent state $|z\rangle$ is defined (Zhang *et al.*, 1990) by the eigenstate of the annihilation operator a with the complex eigenvalue z , i.e.,

$$a|z\rangle = z|z\rangle,$$

and the completeness relation (Ogura and Sekiguchi, 2007)

$$\int \frac{d^2z}{2\pi j} |z\rangle \langle z| = 1.$$

To relate the LCT kernel in quantum mechanics, we use the completeness relation of the coherent state with the arguments of z_1 and z_2 ,

$$\langle Q|U(t)|q\rangle = \int \frac{d^2z_1 d^2z_2}{(2\pi j)^2} \langle Q|z_1\rangle \langle z_1|U(t)|z_2\rangle \langle z_2|q\rangle \quad (5)$$

where

$$\int d^2z \equiv \int d[\text{Re}(z)]d[\text{Im}(z)],$$

and with the aid of the Integration Within Ordered Product (IWOP) technique (Fan and VanderLinde, 1989; Fan *et al.*, 1987) the coherent state representation of the unitary operator is given by

$$\langle z_1|U(t)|z_2\rangle = \frac{1}{\sqrt{s}} \exp \left[-\frac{r}{2s}(z_1^*)^2 + \frac{z_2 z_1^*}{s} + \frac{r^*}{2s}(z_2)^2 - \frac{|z_1|^2}{2} - \frac{|z_2|^2}{2} \right], \quad (6)$$

where (r, s) , $|s^2| - |r^2| = 1$, are related to a classical ray transfer matrix $\begin{pmatrix} a & b \\ c & d \end{pmatrix}$ by

$$s = \frac{1}{2}[a + d + j(b - c)], \quad r = \frac{1}{2}[d - a - j(b + c)], \quad (7)$$

the uni-modularity condition $ad - bc = 1$ is equivalent to $|s^2| - |r^2| = 1$, and the coordinate and momentum representation of the coherent state $|z\rangle$ (Ogura, 2009)

$$\langle Q|z_1\rangle = \frac{1}{\pi^{1/4}} \exp\left[-\frac{Q^2}{2} + \sqrt{2}z_1Q - \frac{z_1^2}{2} - \frac{|z_1|^2}{2}\right], \tag{8}$$

$$\langle q|z_2\rangle = \frac{1}{\pi^{1/4}} \exp\left[-\frac{q^2}{2} - j\sqrt{2}z_2q + \frac{z_2^2}{2} - \frac{|z_2|^2}{2}\right]. \tag{9}$$

By substituting the values of Eqns. (6), (8) and (9) in Eqn. (5), we get

$$\begin{aligned} \langle Q|U(t)|q\rangle &= \int \frac{d^2z_1 d^2z_2}{(2\pi j)^2} \times \frac{1}{\pi^{1/2}} \frac{1}{\sqrt{s}} \\ &\cdot \exp\left[-\frac{Q^2}{2} + \sqrt{2}z_1Q - \frac{z_1^2}{2} - \frac{|z_1|^2}{2}\right] \\ &\cdot \exp\left[-\frac{r}{2s}(z_1^*)^2 + \frac{z_2z_1^*}{s} + \frac{r^*}{2s}(z_2)^2 - \frac{|z_1|^2}{2} - \frac{|z_2|^2}{2}\right] \\ &\cdot \exp\left[-\frac{q^2}{2} + \sqrt{2}z_2^*q - \frac{(z_2^*)^2}{2} - \frac{|z_2|^2}{2}\right]. \end{aligned} \tag{10}$$

This results from solving for z_1 and using the fact (Puri, 2001) that, if $a_1^2 - 4|c_1|^2 > 0$, then

$$\begin{aligned} &\frac{1}{\pi} \int d^2\alpha \exp\left(-a_1|\alpha|^2 + b_1\alpha + b_1^*\alpha^* + c_1\alpha^2 + c_1^*(\alpha^*)^2\right) \\ &= \sqrt{\frac{1}{a_1^2 - 4|c_1|^2}} \exp\left[\frac{b_1^2c_1^* + b_1^{*2}c_1 + a_1|b_1|^2}{a_1^2 - 4|c_1|^2}\right] \end{aligned}$$

and

$$\begin{aligned} &\frac{1}{\pi} \int d^2z_1 \exp\left[-|z_1|^2 + z_1\sqrt{2}Q + \frac{z_2z_1^*}{s} - \frac{z_1^2}{2} - \frac{r}{2s}(z_1^*)^2\right] \\ &= \sqrt{\frac{s}{s-r}} \exp\left[\frac{-Q^2r - \frac{z_2^2}{2s} + \sqrt{2}Qz_2}{s-r}\right]. \end{aligned} \tag{11}$$

Here $a_1 = 1$, $b_1 = \sqrt{2}Q$, $b_1^* = z_2/s$, $c_1 = -1/2$, and $c_1^* = -r/2s$. Similarly, solving for z_2 and substituting the value in Eqn. (10), we finally get

$$\begin{aligned} \langle Q|U(t)|q\rangle &= \sqrt{\frac{1}{\pi(s-s^*-r+r^*)}} \exp\left[\frac{2qQ}{s-s^*-r+r^*}\right. \\ &\left. - \frac{q^2s+s^*-r-r^*}{2s-s^*-r+r^*} - \frac{Q^2s+s^*+r+r^*}{2s-s^*-r+r^*}\right]. \end{aligned} \tag{12}$$

By substituting the value of r and s from Eqn. (7), the LCT kernel is given by

$$\langle Q|U(t)|q\rangle = \sqrt{\frac{1}{j2\pi b}} \exp\left[\frac{j}{2b}(aq^2 - 2qQ + dQ^2)\right]. \tag{13}$$

As a special case of the LCT, when $(a, b, c, d) = (\cos \alpha, \sin \alpha, -\sin \alpha, \cos \alpha)$, Eqn. (13) reduces to the kernel of the fractional Fourier transform (Namias, 1979) $K_\alpha(Q, q)$ as

$$K_\alpha(Q, q) = \sqrt{\frac{1}{2j\pi \sin \alpha}} \exp\left[\frac{j}{2}\left(\frac{q^2 + Q^2}{\cot \alpha} - \frac{2qQ}{\sin \alpha}\right)\right]. \tag{14}$$

Multiplying and dividing Eqn. (13) by $\sqrt{\exp(j\alpha)}$, we finally get

$$L_{(\cos \alpha, \sin \alpha, -\sin \alpha, \cos \alpha)}[x(q)] = \sqrt{\exp(-j\alpha)} F_\alpha[x(q)],$$

where $F_\alpha[x(q)]$ denotes the fractional Fourier transform of $x(q)$. As a special case of the LCT, when $(a, b, c, d) = (0, 1, -1, 0)$, Eqn. (13) reduces to the kernel of Fourier transform as

$$L_{(0,1,-1,0)}[x(q)] = \sqrt{-j} F[x(q)],$$

where $F[x(q)]$ denotes the Fourier transform of $x(q)$. Hence the LCT can extend its utilities and applications and solve some problems that cannot be solved well by the operations defined by Ozaktas *et al.* (2000).

3. Modified identities for the convolution and product theorem

The new definition of the convolution and product theorem for the LCT is as follows.

Theorem 1. (Convolution) *If $F_{(a,b,c,d)}(p)$ is the LCT of $f(x)$ and $G_{(a,b,c,d)}(p)$ is the LCT of $g(x)$, then*

$$\sqrt{j2\pi b} e^{-\frac{j}{2}\left(\frac{d}{b}p^2\right)} F_{(a,b,c,d)}(p) G_{(a,b,c,d)}(p)$$

is a linear canonical transform of $h(x)$, i.e.,

$$\begin{aligned} &L_{(a,b,c,d)}(h(x)) \\ &= \sqrt{j2\pi b} \exp\left[-\frac{j}{2}\left(\frac{d}{b}p^2\right)\right] \\ &\times F_{(a,b,c,d)}(p) G_{(a,b,c,d)}(p), \end{aligned} \tag{15}$$

where

$$h(x) = \int_{-\infty}^{\infty} f(y) g(x-y) \tilde{y}(x,y) dy,$$

is the weighted convolution and the weight function is $\tilde{y}(x,y) = e^{-j\frac{a}{b}y(x-y)}$ while the role of g and f can be interchanged.

Theorem 2. (Product) For any two functions $f(x)$ and $g(x)$, the modified product operation is defined as

$$\begin{aligned}
 m(x) &= g(x) \cdot f(x) \cdot e^{j\frac{a}{2b}x^2} \\
 L_{(a,b,c,d)}(g(x) \cdot f(x) \cdot e^{j\frac{a}{2b}x^2}) \\
 &= G_{(a,b,c,d)}(p) \otimes F_{(a,b,c,d)}(p) \cdot \exp\left[j\frac{d}{b}vp\right]. \quad (16)
 \end{aligned}$$

4. Proposed transform of a convolution by coordinate-momentum representation transform

Using the quantum mechanical notation (Fan and Yue, 2003; Fan *et al.*, 2008), the one-dimensional linear canonical transform of the signal $h(x)$ reads as

$$\begin{aligned}
 L_{(a,b,c,d)}(h(x)) \\
 &= \frac{1}{\sqrt{j2\pi b}} \int_{-\infty}^{\infty} dx \langle p|K|x\rangle \langle x|h\rangle = \langle p|K|h\rangle \quad (17) \\
 &= \langle p|H\rangle = H_{(a,b,c,d)}(p),
 \end{aligned}$$

where we use the notation $\langle x|h\rangle = h(x)$ and $\langle p|K|x\rangle$ gives the representation of the LCT kernel in quantum mechanics. K is named as the LCT operator and

$$\langle x|h\rangle = \int_{-\infty}^{\infty} dy \langle y|f\rangle \langle x-y|g\rangle \exp\left[-j\frac{a}{b}y(x-y)\right]. \quad (18)$$

Substituting the value of Eqn. (18) in (17) results in

$$\begin{aligned}
 L_{(a,b,c,d)}(h(x)) \\
 &= \frac{1}{\sqrt{j2\pi b}} \int_{-\infty}^{\infty} dx \langle p|K|x\rangle \int_{-\infty}^{\infty} dy \langle y|f\rangle \langle x-y|g\rangle \\
 &\quad \cdot \exp\left[-j\frac{a}{b}y(x-y)\right] \quad (19)
 \end{aligned}$$

Rearranging Eqn. (19) we get

$$\begin{aligned}
 L_{(a,b,c,d)}(h(x)) \\
 &= \frac{1}{\sqrt{j2\pi b}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dx dy \langle p|K|x\rangle \langle x-y|g\rangle \langle y|f\rangle \\
 &\quad \cdot \exp\left[-j\frac{a}{b}y(x-y)\right] \quad (20)
 \end{aligned}$$

Substituting $x-y = x'$, i.e., $x = x' + y$ and $y = y$, in Eqn. (20), we have $dx dy = dx' dy$ (Kiusalaas, 2010),

and then replacing x' by x results in

$$\begin{aligned}
 L_{(a,b,c,d)}(h(x)) \\
 &= \frac{1}{\sqrt{j2\pi b}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dx dy \langle p|K|x+y\rangle \langle x|g\rangle \langle y|f\rangle \\
 &\quad \cdot \exp\left[-j\frac{a}{b}yx\right]. \quad (21)
 \end{aligned}$$

Rewriting $\langle p|K|x+y\rangle$ explicitly, we get

$$\begin{aligned}
 \langle p|K|x+y\rangle &= \frac{1}{\sqrt{j2\pi b}} \exp\left[-j\frac{a}{b}yx\right] \\
 &= \frac{1}{\sqrt{j2\pi b}} \\
 &\quad \cdot \exp\left[\frac{jdp^2 - 2jp(x+y) + ja(x+y)^2 - 2jaxy}{2b}\right]. \quad (22)
 \end{aligned}$$

Multiplying and dividing Eqn. (22) by $\exp\left[\frac{j}{2}p^2\right]/\sqrt{j2\pi b}$ result in

$$\begin{aligned}
 &\Rightarrow \frac{\sqrt{j2\pi b}}{j2\pi b} \\
 &\quad \cdot \exp\left[\frac{jdp^2 - 2jpx - 2jpy + jax^2 + jay^2}{2b}\right] \\
 &\quad \cdot \exp\frac{jdp^2 - jdp^2}{2b} \quad (23)
 \end{aligned}$$

Rearranging Eqn. (23) we get

$$\begin{aligned}
 &\Rightarrow \frac{\sqrt{j2\pi b}}{j2\pi b} \left[\exp\left\{\frac{jdp^2 - 2jpx + jax^2}{2b}\right\} \right. \\
 &\quad \left. \cdot \exp\left\{\frac{jdp^2 - 2jpy + jay^2}{2b}\right\} \exp\left\{\frac{-jdp^2}{2b}\right\} \right]. \quad (24)
 \end{aligned}$$

Substituting Eqn. (24) in (21) results in

$$\begin{aligned}
 L_{(a,b,c,d)}(h(x)) \\
 &= \frac{\sqrt{j2\pi b}}{j2\pi b} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dx dy \langle x|g\rangle \langle y|f\rangle \\
 &\quad \cdot \left[\exp\left\{\frac{jdp^2 - 2jpx + jax^2}{2b}\right\} \right. \\
 &\quad \left. \cdot \exp\left\{\frac{jdp^2 - 2jpy + jay^2}{2b}\right\} \cdot \exp\left\{\frac{-jdp^2}{2b}\right\} \right]. \quad (25)
 \end{aligned}$$

Rewriting Eqn. (25) results in

$$\begin{aligned}
 L_{(a,b,c,d)}(h(x)) \\
 &= \frac{\sqrt{j2\pi b}}{j2\pi b} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dx dy \langle p|K|x\rangle \langle x|g\rangle \\
 &\quad \cdot \langle p|K|y\rangle \langle y|f\rangle \exp\left[\frac{-jdp^2}{2b}\right]. \quad (26)
 \end{aligned}$$

Rearranging Eqn. (26), we get

$$L_{(a,b,c,d)}(h(x)) = \sqrt{j2\pi b} \left[\frac{1}{\sqrt{j2\pi b}} \int_{-\infty}^{\infty} dx \langle p|K|x\rangle \langle x|g\rangle \cdot \frac{1}{\sqrt{j2\pi b}} \int_{-\infty}^{\infty} dy \langle p|K|y\rangle \langle y|f\rangle \cdot \exp\left\{\frac{-jdp^2}{2b}\right\} \right]. \quad (27)$$

Rewriting Eqn. (27) results in

$$L_{(a,b,c,d)}(h(x)) = \sqrt{j2\pi b} e^{-\frac{j}{2}\left(\frac{a}{b}p^2\right)} F_{(a,b,c,d)}(p) G_{(a,b,c,d)}(p). \quad (28)$$

This is just a new convolution theorem under the LCT, derived by representation transformation in quantum mechanics. Hence Eqn. (15) is obtained and the theorem is proved.

The reciprocal transform of Eqn. (28) can be obtained by writing the definition of the inverse LCT and is given by

$$L^{-1}_{(a,b,c,d)}(H(p)) = \int_{-\infty}^{\infty} dp \langle p|K|x\rangle^* \cdot H_{(a,b,c,d)}(p), \quad (29)$$

where * indicates the complex conjugate. By using the theory of representation transformation in quantum mechanics, Eqn. (29) results in

$$L^{-1}_{(a,b,c,d)}(H(p)) = \int_{-\infty}^{\infty} dp \langle x|K^\dagger|p\rangle \cdot \langle p|H\rangle = \int_{-\infty}^{\infty} dp \langle x|K^\dagger K|h\rangle = \langle x|h\rangle = h(x). \quad (30)$$

Seeing FT as a special case of the LCT, when $(a, b, c, d) = (0, 1, -1, 0)$, Eqn. (28) becomes

$$L_{(0,1,-1,0)}(h(x)) = \sqrt{j2\pi} F_{(0,1,-1,0)}(p) G_{(0,1,-1,0)}(p). \quad (31)$$

Similarly, seeing the FRFT as a special case of the LCT, when $(a, b, c, d) = (\cos \alpha, \sin \alpha, -\sin \alpha, \cos \alpha)$, Eqn. (28) becomes

$$F_\alpha(h(x)) = \sqrt{j2\pi \sin \alpha} e^{-\frac{j}{2}p^2 \cot \alpha} F_\alpha(p) G_\alpha(p), \quad (32)$$

and the FT as a special case of the FRFT, $\alpha = \pi/2$, (32) becomes

$$F_{\pi/2}(h(x)) = \sqrt{j2\pi} F_{\pi/2}(p) G_{\pi/2}(p). \quad (33)$$

Equations (31) and (32) are special cases of the LCT.

5. Proposed transform of a product by coordinate-momentum representation transform

The one-dimensional linear canonical transform of $m(x)$ in the context of quantum mechanics is given by

$$L_{(a,b,c,d)}(m(x)) = \frac{1}{\sqrt{j2\pi b}} \int_{-\infty}^{\infty} dx \langle p|K|x\rangle \langle x|m\rangle. \quad (34)$$

From the definition of the modified product theorem as given by Eqn. (16),

$$L_{(a,b,c,d)}(m(x)) = \frac{1}{\sqrt{j2\pi b}} \int_{-\infty}^{\infty} dx \langle x|f\rangle \langle x|g\rangle \langle p|K|x\rangle \cdot \exp\left[j\frac{a}{2b}x^2\right], \quad (35)$$

the one-dimensional inverse linear canonical transform in the context of quantum mechanics is given by

$$\langle x|g\rangle = \sqrt{\frac{-1}{j2\pi b}} \int_{-\infty}^{\infty} dv \langle v|G\rangle \langle v|K|x\rangle^*, \quad (36)$$

where * indicates the complex conjugate. Rewriting Eqn. (36) results in

$$\langle x|g\rangle = \sqrt{\frac{-1}{j2\pi b}} \int_{-\infty}^{\infty} dv \langle v|G\rangle \langle x|K^\dagger|v\rangle. \quad (37)$$

Substituting the value of Eqn. (37) in (35) results in

$$L_{(a,b,c,d)}(m(x)) = \left|\frac{1}{2\pi b}\right| \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dv dx \langle v|G\rangle \cdot \langle x|K^\dagger|v\rangle \langle x|f\rangle \cdot \langle p|K|x\rangle \exp\left[j\frac{a}{2b}x^2\right]. \quad (38)$$

Solving for $\langle x|K^\dagger|v\rangle \cdot \langle p|K|x\rangle$ results in

$$\Rightarrow \exp\left[\frac{-jdv^2 + 2jvpx - jax^2}{2b}\right] \cdot \exp\left[\frac{jdp^2 - 2jpx + jax^2}{2b}\right]. \quad (39)$$

Substituting the value of Eqn. (39) in (38) results in

$$\begin{aligned}
 &L_{(a,b,c,d)}(m(x)) \\
 &= \left| \frac{1}{2\pi b} \right| \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dv dx \langle v|G\rangle \langle x|f\rangle \cdot \exp \left[j \frac{a}{2b} x^2 \right] \\
 &\cdot \exp \left[\frac{-jdv^2 + 2jvx - jax^2}{2b} \right] \\
 &\cdot \exp \left[\frac{jdp^2 - 2jpx + jax^2}{2b} \right].
 \end{aligned} \tag{40}$$

Multiplying and dividing Eqn. (40) by $\exp \left[\frac{-j}{2} \frac{d}{b} v(v - 2p) \right]$ results in

$$\begin{aligned}
 &L_{(a,b,c,d)}(m(x)) \\
 &= \left| \frac{1}{2\pi b} \right| \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dv dx \langle v|G\rangle \langle x|f\rangle \\
 &\cdot \exp \left[\frac{-j}{2} \left\{ \frac{dv^2 - 2vx - dp^2 + 2px - ax^2}{b} \right. \right. \\
 &\left. \left. + \frac{dv^2 - 2dpv - dv^2 + 2dpv}{b} \right\} \right].
 \end{aligned} \tag{41}$$

Rewriting Eqn. (41), we get

$$\begin{aligned}
 &L_{(a,b,c,d)}(m(x)) \\
 &= \left| \frac{1}{2\pi b} \right| \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dv dx \langle v|G\rangle \langle x|f\rangle \\
 &\cdot \exp \left[\frac{j}{2} \left\{ \frac{d}{b} (p^2 + v^2 - 2pv) - \frac{2x}{b} (p - v) + \frac{ax^2}{b} \right\} \right] \\
 &\cdot \exp \left[-j \frac{d}{b} v^2 + j \frac{d}{b} pv \right].
 \end{aligned} \tag{42}$$

Rewriting Eqn. (42) results in

$$\begin{aligned}
 &L_{(a,b,c,d)}(m(x)) \\
 &= \left| \frac{1}{2\pi b} \right| \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dv dx \langle v|G\rangle \langle x|f\rangle \\
 &\cdot \exp \left[\frac{j}{2} \left\{ \frac{d}{b} (p - v)^2 - \frac{2x}{b} (p - v) + \frac{ax^2}{b} \right\} \right] \\
 &\cdot \exp \left[j \frac{d}{b} v(p - v) \right].
 \end{aligned} \tag{43}$$

From the representation theory of quantum mechanics,

Eqn. (43) can be written as

$$\begin{aligned}
 &L_{(a,b,c,d)}(m(x)) \\
 &= \sqrt{\frac{-1}{j2\pi b}} \int_{-\infty}^{\infty} dv \langle v|G\rangle \langle p - v|f\rangle \\
 &\cdot \exp \left[j \frac{d}{b} v(p - v) \right] \\
 &= G_{(a,b,c,d)}(p) \otimes F_{(a,b,c,d)}(p) \\
 &\cdot \exp \left[j \frac{d}{b} vp \right].
 \end{aligned} \tag{44}$$

This is just a new product theorem under the LCT, derived by representation in quantum mechanics. Hence Eqn. (16) is solved and the theorem is proved.

6. Simulation results

The convolution theorem in the LCT domain given by Deng *et al.* (2006) and Wei *et al.* (2012; 2009) is compared with the proposed convolution theorem by simulating on system with an Intel core™ i3-330M 2.13 GHz processor with 3 GB RAM. The convolution operation of a rectangular function $x(t)$ of unit amplitude is performed with itself, i.e., $(x \otimes x)(t)$. As a result of the convolution operation, one obtains a triangular (Bartlett) function with a duration double of that of the rectangular function as shown in Fig. 1. Then the LCT of convolution operations defined by Deng *et al.* (2006) and Wei *et al.* (2012; 2009) is compared with the LCT of the proposed convolution operation for $(a, b, c, d) = (0.707, 0.707, -0.707, 0.707)$ and $(a, b, c, d) = (0.5, 0.866, -0.866, 0.5)$ as shown in Figs. 2 and 3, respectively. Simultaneously, the LCT of the triangular function is also evaluated for the same values of (a, b, c, d) to make a comparison. It is shown that the real (Re), imaginary (Im) and absolute (Abs) components of the proposed convolution theorem for the LCT resemble maximally the different components of the LCT of a triangular function for $(a, b, c, d) = (0.707, 0.707, -0.707, 0.707)$ and $(a, b, c, d) = (0.5, 0.866, -0.866, 0.5)$.

7. Application

To illustrate the proposed work, an example of filtering is presented. For filtering applications, the time-frequency plot of the original signal and the corrupted signal plays an important role. Many time-frequency signal representations are available, such as the Wigner Distribution (WD) (Wigner, 1932; Hlawatsch and Boudreaux-Bartels, 1992; Classen and Mecklenbrauker, 1980), the Short Time Fourier Transform (STFT) (Hlawatsch and Boudreaux-Bartels, 1992; Portnoff, 1980; Cohen, 1989), the ambiguity function (Hlawatsch

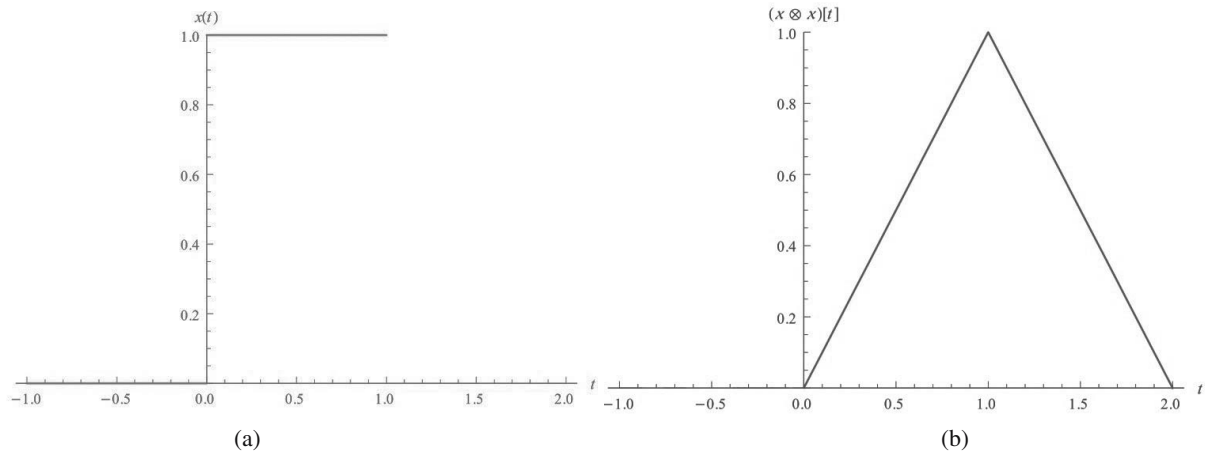


Fig. 1. Rectangular function $x(t)$ (a), convolved signal $(x \otimes x)(t)$, i.e., triangular (Bartlett) window (b).

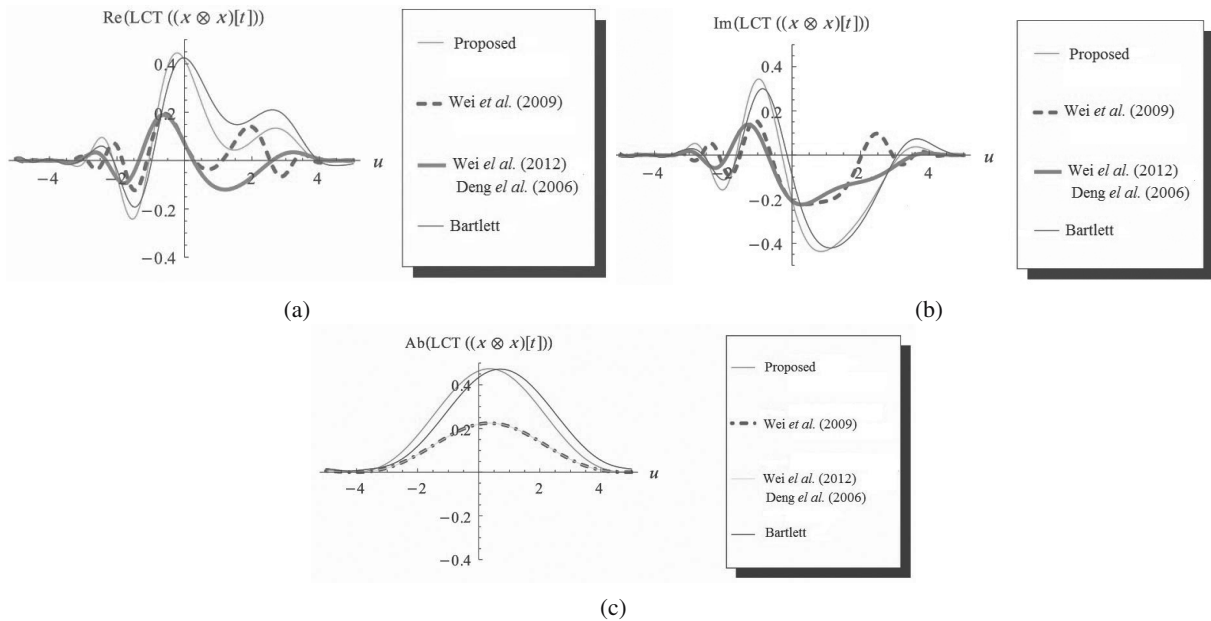


Fig. 2. Real value (a), imaginary value (b), absolute value of the LCT of $(x \otimes x)(t)$ by using the Deng *et al.* (2006) and Wei *et al.* (2012; 2009), Bartlett and proposed methods for $(a, b, c, d) = (0.707, 0.707, -0.707, 0.707)$ (c).

and Boudreaux-Bartels, 1992; Cohen, 1989), the Gabor transform (Swiercz, 2010) and the spectrogram (Hlawatsch and Boudreaux-Bartels, 1992; Portnoff, 1980; Cohen, 1989). Out of these techniques, the WD is a powerful (Bouachache and Rodriguez, 1984) time-frequency analysis tool and is used in this paper. The concept of the canonical filter design is understood from the definition of canonical convolution (Pei and Deng, 2001), given by

$$g(t) = L_{(d,-b,-c,a)} \left(L_{(a,b,c,d)}(h(t)) \cdot L_{(a,b,c,d)}(x(t)) \right),$$

where

$$H(u) = L_{(a,b,c,d)}(h(t))$$

represents the transfer function of the pass-stop band canonical filter and

$$L_{(d,-b,-c,a)} \left(L_{(a,b,c,d)}(f(t)) \right) = f(t)$$

gives the reversibility property of the LCT. Consider an original signal

$$x(t) = 2 \exp \left[-\frac{1}{8}(t+1)^2 - j0.5t^2 \right].$$

Let this signal be corrupted by an Additive White Gaussian Noise (AWGN) of 5 dB Signal to Noise Ratio

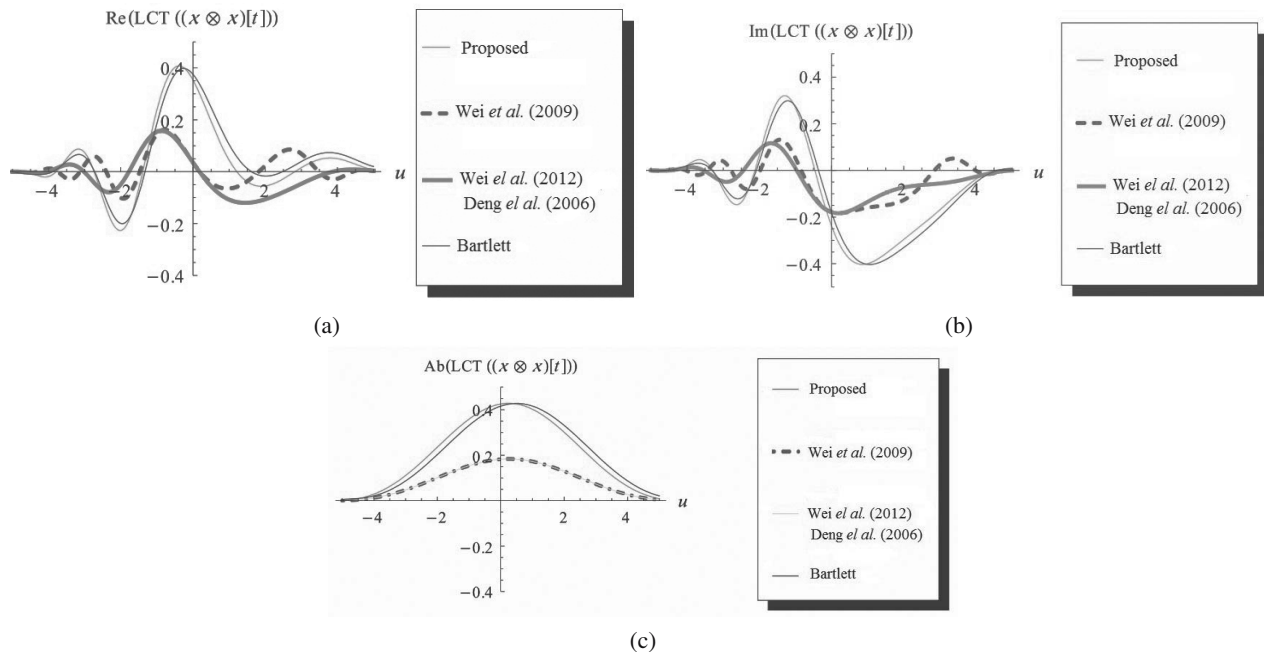


Fig. 3. Real value (a), imaginary value (b), absolute value of the LCT of $(x \otimes x)(t)$ by using the Deng *et al.* (2006) and Wei *et al.* (2012; 2009), Bartlett and proposed methods for $(a, b, c, d) = (0.5, 0.866, -0.866, 0.5)$ (c).

(SNR). As shown in Fig. 4(a), the original signal $x(t)$ is plotted. Figure 4(c) shows the corrupted signal. Figures 4(b) and 4(d) show the WD of the original signal and corrupted signal, respectively. Following the method given by Pei and Deng (2001), we found the optimal filtering domain is the canonical domain with $(a, b, c, d) = (0.16, 1, -1, 0)$. Figures 4(e) and 4(f) show a comparison between the imaginary and real part of the exact signal and the proposed signal output of the filter with the Mean Square Error (MSE) equal to 0.00049. As a comparison with the convolution theorems already derived, the number of chirp functions required by the proposed convolution process is two (cf. Eqn. (18)) and seven (cf. Eqn. (24)), compared with seven and six chirp functions by Wei *et al.* (2009) or three and seven chirp functions by Deng *et al.* (2006) for the left and right-hand sides of the convolution process, respectively. Hence the proposed convolution theorem also reduces the computational complexity.

8. Conclusion

We defined the convolution and product operations of two functions in the LCT domain and derived the convolution and product theorem for linear canonical transformation by employing the coordinate-momentum representation and the LCT operator. The derivation is direct and concise since we used Dirac's representation theory (Huang *et al.*, 2006). It is seen from the results of Figs. 2 and 3 that the proposed weighted convolution theorem gives

better results than the convolution theorem by Deng *et al.* (2006) and Wei *et al.* (2012; 2009), as the results given by the proposed theorem resemble maximally the shape of the LCT of the triangular function. Further, an application of filtering was presented with the proposed convolution theorem and it was found that with the help of proposed theorem the signal is recovered with a minimum mean square error. The derived approach can be used for defining the canonical correlation theorem for the LCT, which will play an important role in signal spectrum analysis, beamforming and signal detection in the canonical domain.

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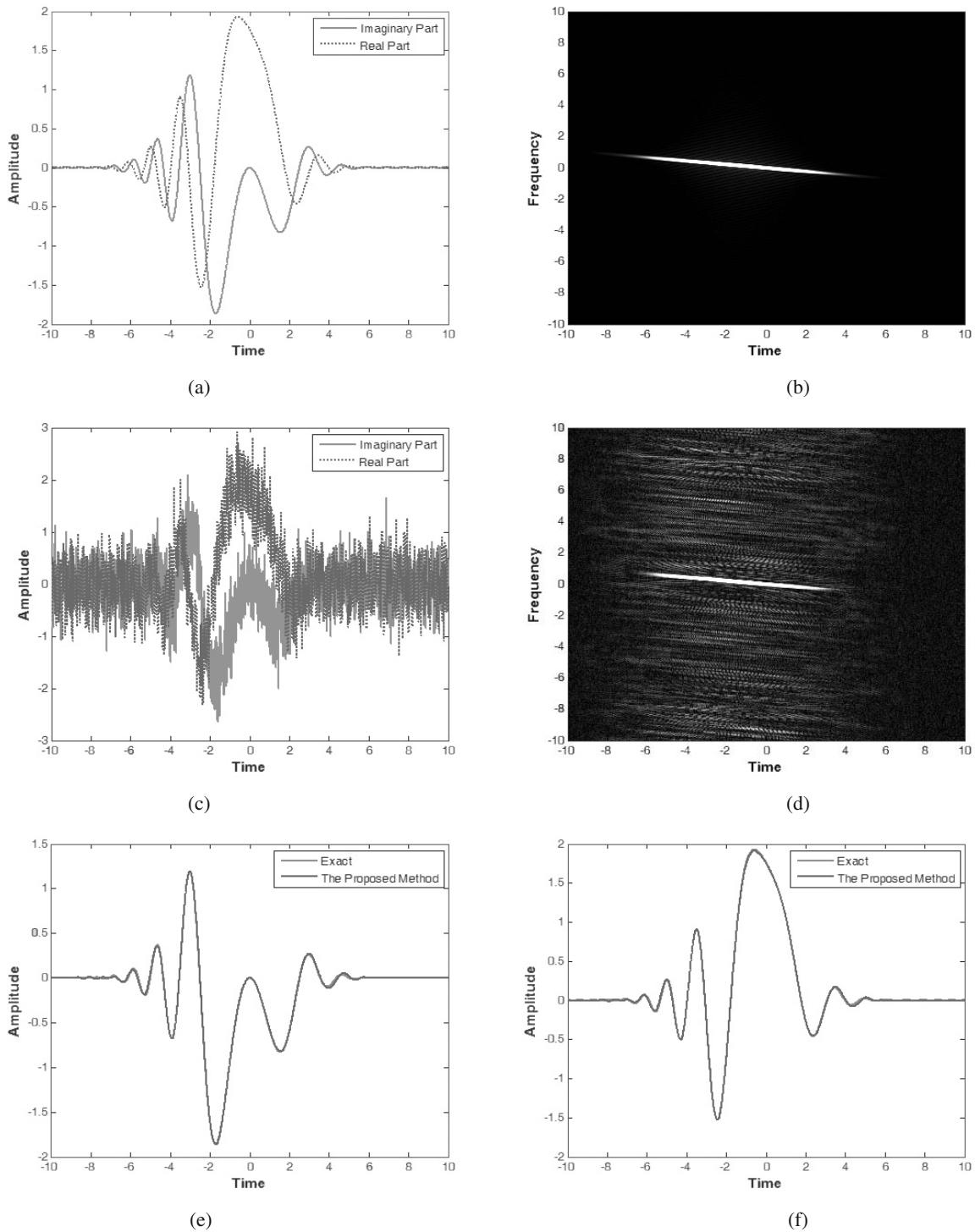


Fig. 4. Original signal (a), WD of the original signal (b), corrupted signal (c), WD of the corrupted signal (d), imaginary part: comparison of the recovered and original signals (e), real part: comparison of the recovered and original signals (f).

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